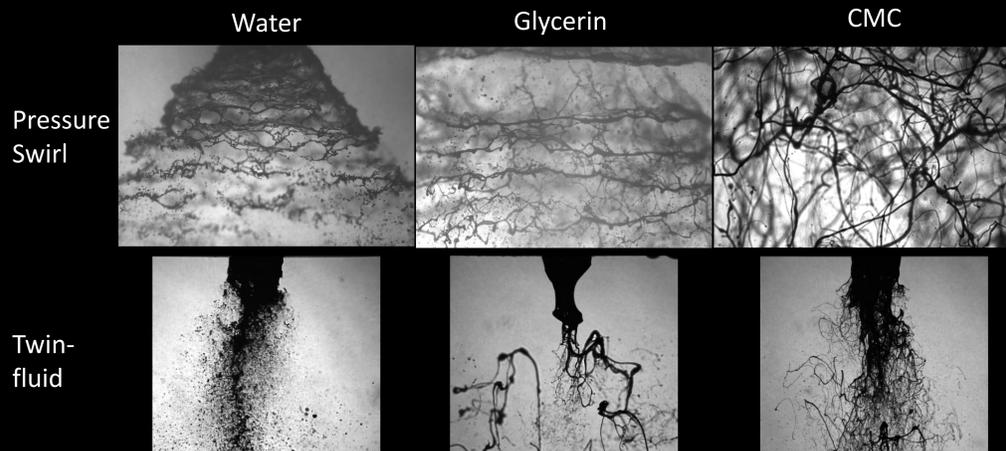


Spray Characterization at Industrially Relevant Conditions

1. Background and Introduction

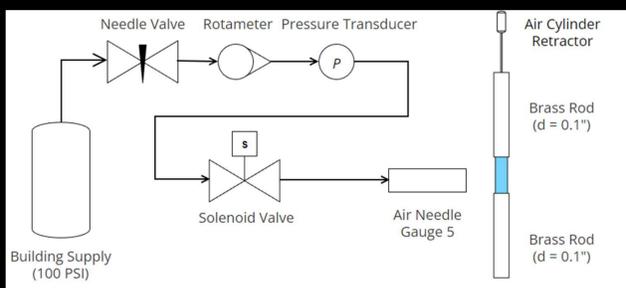
- We have modelled atomization of inviscid and Newtonian fluid in twin fluid and swirl nozzles in previous term.
- Ligaments of highly viscous, polymeric fluid stretch and become thinner before they break up into droplets
- We need to study the stretching and breakup mechanism of ligaments for these complex fluid to modify our models.



2. Objectives

- Breakup of ligaments of fluids in cross flows.
- Spray droplet size measurement of polymeric fluids and slurries.
- Models for the atomization of high viscosity/polymeric fluids and slurries in swirl and twin fluid nozzles based on ligament dynamics.

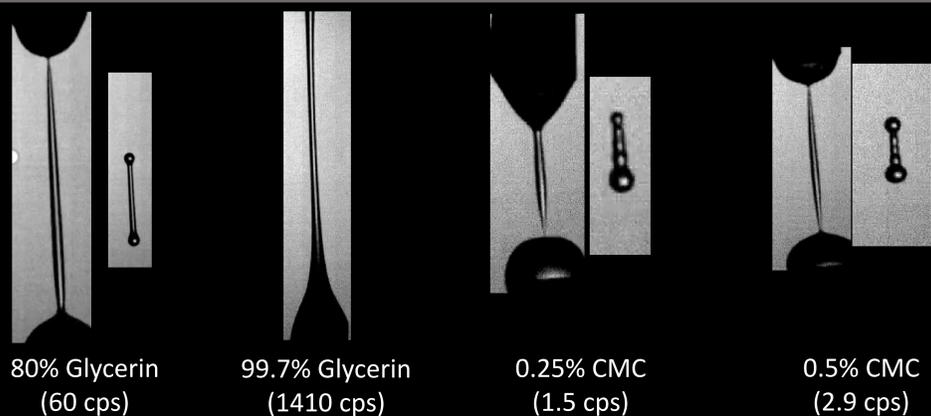
3. Experiment setup



Retractor speed: 0.4m/s Air velocity: 30-70 m/s

Fluid :	High Viscous Newtonian fluid	80 % Glycerin	99.7% Glycerin
	Polymeric fluid	0.25% CMC	0.5% CMC
	Solvent and slurry	Neodol	Slurry with Neodol

4. Pulling without air cross flow



6. Summary

- Designed an experiment to study the breakup of liquid ligaments in cross flows.
- Initial experiments show
 - ligament stretching and thinning, followed by bead and then droplet formation or
 - bag formation, followed by bag breakup and ligament formation by the rims of the bag, followed by ligaments stretching and breakup.
 - Droplet sizes appear smaller in the bag-ligament-drop breakup mechanism.
- Modified a model to describe the flapping process

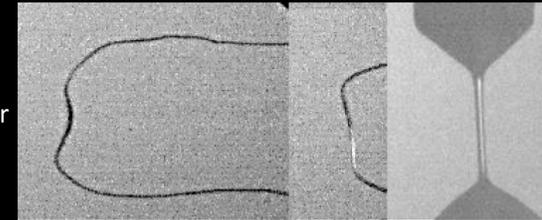
7. Future work

- More detailed ligament breakup experiments and measurement of droplet sizes after breakup, for polymeric, and slurry ligaments.
- Identify effect of elongational viscosity on breakup mechanism and droplet size.

5. Pulling with air crossflow

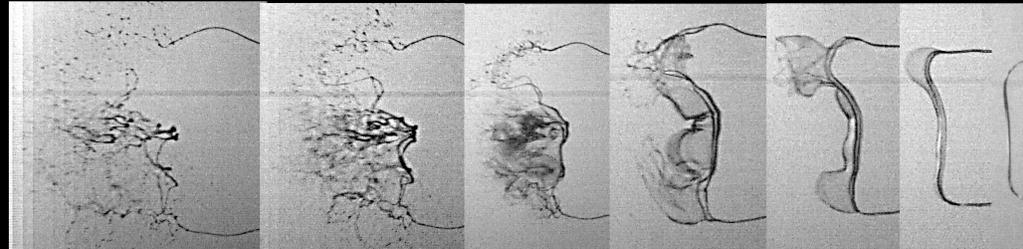
The cylindrical ligament is initially deformed due to air stream.

80% Gly-water, $u_g = 50m/s$, $d_{lig} = 210\mu m$, $We = 10$

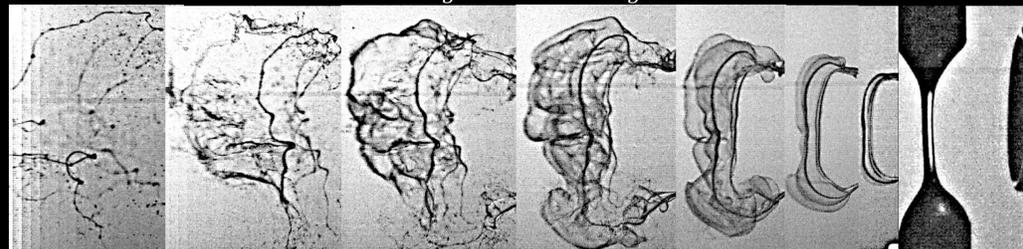


Then depending on the Weber number $We = \frac{\rho_g U_g^2 d_{lig,0}}{\sigma}$, the ligament will either experience stretching or bag breakup and $We_c = 11$.

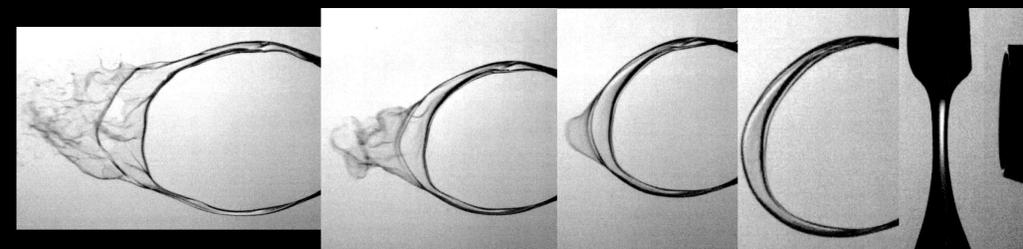
80% Gly-water, $u_g = 72m/s$, $d_{lig} = 300\mu m$, $We = 28$



80% Gly-water, $u_g = 72m/s$, $d_{lig} = 480\mu m$, $We = 46$

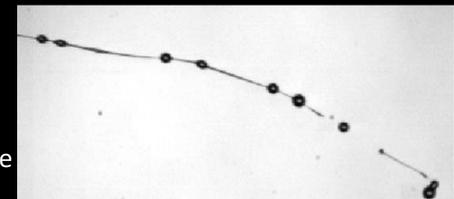


99.7% Glycerin, $u_g = 72m/s$, $d_{lig} = 570\mu m$, $We = 54$

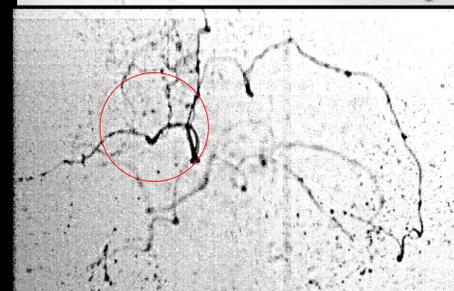


There are two mechanism of bead formation.

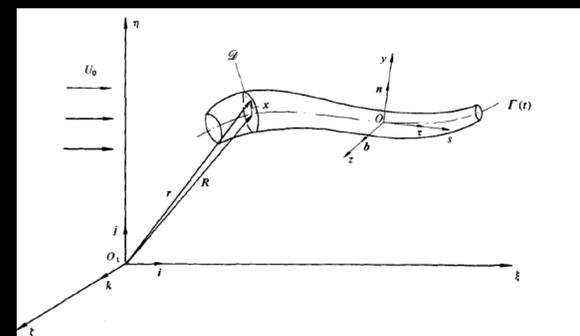
1- Classical Rayleigh-Plateau Instability – In this mechanism, small disturbances on a cylindrical liquid ligament grow, causing the breakup of the ligament. The bead formation appear when the thinning rate of the ligament is smaller than the growth rate of the instability.



2- Undulations formed by a flapping ligament. Ligaments form bead as they go through large sinusoidal oscillation.



- We modified a model by proposed by Yarin A.L. (1993), which provides a set of equations for
 - the movement of ligament central axis $\Gamma(r)$
 - the movement of a fluid particle in the cross section centered at some point $r = R$



$$\frac{\partial \lambda a^2}{\partial t} + \frac{\partial W a^2}{\partial s} = 0 \quad (\text{mass conservation})$$

$$\frac{\partial \lambda a^2 v_n}{\partial t} + V_t \frac{a^2}{\lambda} \frac{\partial v_n \lambda}{\partial s} + \frac{\partial W v_n a^2}{\partial s} + a^2 W V_t \lambda k = \frac{1}{Re} \left(\frac{\partial Q_n}{\partial s} + \lambda P k \right) - J \frac{k a^2}{\lambda} \quad (\text{momentum})$$

$$v = V + \Omega \times x + \delta x + v' \quad (\text{affine transformation in cross section})$$

Compared to classical Rayleigh-Plateau instability, this model describes the ligament dynamics based on

- deforming coordinate system fixed on the ligament central axis, and
- changing ligament cross section.

